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Investigation of the accuracy of an approximate solution of the initial boundary value problem of thermal conductivity with boundary conditions of the third kind

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Abstract: The paper compares two approximations of the solution of the initial boundary value problem of thermal conductivity with boundary conditions of the third kind for the case of a temperature shock of a homogeneous thin plate. A second-order approximation with respect to the plate thickness is obtained. The results of the numerical simulation are compared with the known first-order approximation with respect to the plate thickness. The results obtained demonstrate good convergence. They can be used in controlling the motion of a small spacecraft with elastic structural elements. It is necessary to level out the effect of temperature shock disturbances on the angular motion of a small spacecraft.

Keywords: thermal conductivity problem; boundary conditions of the third kind; temperature shock; thin plate.

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1 Introduction

From a theoretical point of view, the one-dimensional problem of thermal conductivity is well studied. So, back in 1905, work Einstein (1905) on the thermal motion of suspended particles was published. However, interest in this task has not waned so far. In Muftu (2022), a finite element form was developed for solving one-dimensional equations of thermal conductivity in a steady state with simple boundary conditions. Paper Biswas et al. (2023) is devoted to the search for a new nonlinear shape between two nodal points of a finite element model of one-dimensional thermal conductivity. In Sedelnikov and Orlov (2020), the law of control of a small spacecraft was developed taking into account the temperature effects on it from the Sun in the framework of a one-dimensional thermal conductivity problem. Niu et al. (2021) consider traditional tools for analysing thermal conductivity in reactor systems to be ineffective. In their opinion, these approaches face serious difficulties in developing models, have low accuracy and poor convergence. In Niu et al. (2021), an automatic differentiation method is presented, which allows for automatic numerical calculation of derived functions when solving a one-dimensional thermal conductivity problem. It was about classical problems in the correct formulation according to Hadamard and Morse (1953).

Incorrect tasks according to Hadamard and Morse (1953), first of all, inverse tasks are unlikely to ever lose relevance. It is no coincidence that many works today are devoted to solving incorrect problems. In Jahangiri et al. (2019), a numerical method was used to solve the one-dimensional inverse problem of thermal conductivity, which is a

combination of the Haar wavelet collocation (Mehandiratta et al., 2020) and the Tikhonov regularisation approach (Calvetti and Reichel, 2003). In Molhem and Pourgholi (2008), a numerical algorithm is presented for solving the one-dimensional inverse problem of thermal conductivity in a dimensionless form. The algorithm combines the use of the finite difference method with the solution of an ordinary differential equation. In Pourgholi et al. (2013), a numerical method for solving the inverse problem of thermal conductivity based on the Sinc-Galerkin method is proposed. A stable numerical solution to this problem is obtained. It is noted in Wang et al. (2024) that inverse problems of thermal conductivity play an important role in the engineering field. A method of fast Bayesian parallel sampling is proposed for large-scale estimation of the parameters of the reference three-dimensional inverse problem of thermal conductivity. This method allows you to quickly calculate the parameters on a scale of 10^5 by reducing the dimension of the problem domain, depending on space and time.

Boundary conditions of the third kind should be mentioned separately (Eremin and Gubareva, 2019). They are essentially nonlinear (Sidorov, 1969). This makes it more difficult to obtain an approximate analytical solution to the problem (Sedelnikov et al., 2021).

Even more difficult is the formulation of the thermal conductivity problem for rapidly alternating processes (Choi et al., 2022). One of these processes is temperature shock (Kartashov, 2012). In practice, a temperature shock occurs when a spacecraft with large solar panels is immersed in the shadow and comes out of the shadow (Sedelnikov et al., 2023). In this case, thermal vibrations are excited (Johnston and Thornton, 2000), which can lead to a deterioration in the quality of the spacecraft's solution to its target tasks (Serdakova, 2023). First of all, this applies to Earth remote sensing spacecraft (Khnyryova, 2023). However, when using thin ROSA solar panels (Lee et al., 2013), thermal fluctuations can significantly affect the dynamics of the spacecraft (Shen et al., 2017). During the tests of these panels on the International space Station, thermal fluctuations didn't allow them to collapse (Chamberlain et al., 2021). This shows that a temperature shock can affect the controllability of a spacecraft (Liu et al., 2022). Especially when it comes to small spacecraft (Liu et al., 2019). In Sedelnikov et al. (2024b), new technologies for controlling a small spacecraft are considered, taking into account disturbances from a temperature shock. It is noted that in order to construct an effective control law for a small spacecraft, it is necessary to have approximate analytical dependences of perturbations on a temperature shock. The high load of the on-board computer with the target tasks (Prokopyev et al., 2020) does not allow the installation of resource-intensive applications for the numerical evaluation of these disturbances (Joseph et al., 2018). Therefore, the derivation of approximate analytical dependencies is an important practical task for effective motion control of a small spacecraft.

The purpose of this work is to study the accuracy of the first approximation of solving a one-dimensional thermal conductivity problem with boundary conditions of the third kind for a thin plate under the action of a temperature shock. The first approximation is described in Sedelnikov et al. (2024a) and demonstrates good convergence with the numerical solution in the Ansys package.

2 Setting the task

Let consider the initial boundary value problem of one-dimensional thermal conductivity in the form (Sedelnikov et al., 2021, 2024a, 2024b):

$$\begin{cases} \frac{\partial T(z,t)}{\partial t} = a \frac{\partial^2 T(z,t)}{\partial z^2}, 0 \leq z \leq h, t > 0; \\ \left(\lambda \frac{\partial T(h,t)}{\partial z} \right) = Q - e\Theta (T^4(h,t) - T_c^4), z = h, t > 0; \\ \left(\lambda \frac{\partial T(0,t)}{\partial z} \right) = -e\Theta (T^4(0,t) - T_c^4), z = 0, t > 0; \\ T(z,0) = T_0 = \text{const}, 0 \leq z \leq h, t = 0, \end{cases}$$

where

a thermal conductivity coefficient of the material

h plate thickness

λ coefficient of thermal conductivity

e the degree of blackness of the plate material

Θ the Stefan-Boltzmann constant

T_c ambient temperature of the plate

Q heat flow incident on the surface of the plate.

Neglecting the ambient temperature (approximately 3 K (Guo et al., 2024)) in comparison with the temperature of the plate, we rewrite the system of equations as:

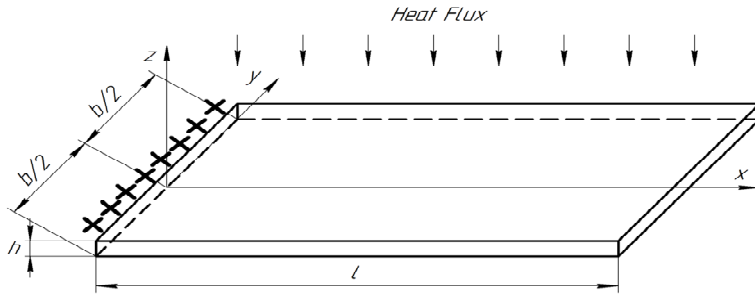
$$\begin{cases} \frac{\partial T(z,t)}{\partial t} = a \frac{\partial^2 T(z,t)}{\partial z^2}, 0 \leq z \leq h, t > 0; \\ \left(\lambda \frac{\partial T(h,t)}{\partial z} \right) = Q - e\Theta T^4(h,t), z = h, t > 0; \\ \left(\lambda \frac{\partial T(0,t)}{\partial z} \right) = -e\Theta T^4(0,t), z = 0, t > 0; \\ T(z,0) = T_0 = \text{const}, 0 \leq z \leq h, t = 0. \end{cases} \quad (1)$$

A homogeneous thin rectangular plate with constant thermophysical properties over the entire temperature range is considered as an object of research. The plate is rigidly fixed with one edge. The other three edges of the plate are free (Figure 1). Heat losses through the side surface of the plate are considered negligible. At the initial moment of time, the temperature field of the plate was homogeneous. At the moment of the thermal shock, the plate had a flat shape. The Q flow is perpendicular to the plate plane at all times.

Such a setting in Sedelnikov and Orlov (2020) and Sedelnikov et al. (2021) is defined as the most dangerous in terms of the effect of temperature shock on the dynamics of a small spacecraft. The flow perpendicular to the plate ensures maximum heating of its surface layer. In this case, the disturbances from the temperature shock will be maximal (Sedelnikov et al., 2024b).

Thus, the problem statement under consideration is relevant from the point of view of estimating the maximum effect of a temperature shock. Based on it, it is possible to make a decision on taking into account the temperature shock when modelling the motion of a small spacecraft. Therefore, the studies conducted in this paper on the accuracy of the approximate solution of the initial boundary value problem (1), presented in Sedelnikov et al. (2024a), are relevant and important.

Figure 1 The appearance of the plate being modelled in operation



3 Output of an approximate solution

For the output, we use the expansion (Sedelnikov et al., 2024a; Korneyev et al., 2019):

$$T_n(z, t) = \sum_{i=0}^n \xi_i(t) z^i \quad (2)$$

It is a well-known Fourier decomposition of a function of two variables (Watts, 2012). There are studies on the incorrectness of using such decomposition to describe complex processes (Arinchev, 2022). However, in this paper it is considered suitable for obtaining an approximate analytical dependence for the temperature field of the plate in the framework of the one-dimensional problem of thermal conductivity. This can be a problem for solving two-dimensional (Sedelnikov et al., 2021b) or three-dimensional thermal conductivity problems (Cheng and Ge, 2010).

As in Sedelnikov et al. (2024a), we take into account the first four terms in expansion (2):

$$T_3(z, t) = \xi_0(t) + \xi_1(t)z + \xi_2(t)z^2 + \xi_3(t)z^3 \quad (3)$$

Taking into account the derivatives, substituting (3) into the first equation (1) gives the following equation Sedelnikov et al. (2025):

$$\frac{d\xi_0(t)}{dt} + \frac{d\xi_1(t)}{dt}z + \frac{d\xi_2(t)}{dt}z^2 + \frac{d\xi_3(t)}{dt}z^3 = 2a[\xi_2(t) + 3\xi_3(t)z] \quad (4)$$

However, in Sedelnikov et al. (2024a), equation (5) was linearised with respect to the variable z . The plate in question is thin. The variable z varies from 0 to h . Therefore, such linearisation is fully justified. In this paper, it is proposed to take into account the terms of the second order relative to z in equation (4). After that, compare the results obtained

with the linearised equation (4) and draw conclusions about the correctness and adequacy of the linearised equation (4).

Let's group the terms in (4), neglecting the terms above the second order with respect to z :

$$\left[\frac{d\zeta_2(t)}{dt} \right] z^2 + \left[\frac{d\zeta_1(t)}{dt} - 6a\zeta_3(t) \right] z + \left[\frac{d\zeta_0(t)}{dt} - 2a\zeta_2(t) \right] = 0 \quad (5)$$

The functions in square brackets of equation (5) are generally independent. Therefore, in order to satisfy equation (5), we require that all of them be equal to zero at the same time:

$$\begin{cases} \frac{d\zeta_2(t)}{dt} = 0; \\ \frac{d\zeta_1(t)}{dt} - 6a\zeta_3(t) = 0; \\ \frac{d\zeta_0(t)}{dt} - 2a\zeta_2(t) = 0 \end{cases} \quad (6)$$

Conditions (6) are in some sense similar to the conditions in the Rayleigh-Ritz method, where each term of the expansion must separately satisfy the boundary conditions (Abdallah and Madjid, 2024). Let's arrange the equations in system (6) by terms of expansion and take the obvious integrals. Then (6) transforms to the form:

$$\begin{cases} \zeta_0(t) = a\beta_1 t^2 + \beta_2; \\ \frac{d\zeta_1(t)}{dt} = 6a\zeta_3(t); \\ \zeta_2(t) = \beta_1 t, \end{cases} \quad (7)$$

where β_1 и β_2 – integration constants.

One of the integration constants can be found from the initial conditions. The fourth equation of the system of equations (1) for $z = 0$ will take the form: $T(0, 0) = T_0$. According to decomposition (3), for $z = 0$ we have:

$$T_3(0, t) = \zeta_0(t). \quad (8)$$

At $t = 0$ get:

$$T_3(0, 0) = \zeta_0(0) = T_0. \quad (9)$$

Substitute $t = 0$ into the first equation of the system of equations (7) and obtain:

$$\beta_2 = \zeta_0(0) = T_0. \quad (10)$$

Then the system of equations (7) will take the final form:

$$\begin{cases} \zeta_0(t) = a\beta_1 t^2 + T_0; \\ \frac{d\zeta_1(t)}{dt} = 6a\zeta_3(t); \\ \zeta_2(t) = \beta_1 t. \end{cases} \quad (11)$$

If the decomposition functions (3) satisfy the conditions (11), then the decomposition (3) itself will satisfy the equation of thermal conductivity. In the system of equations (11), only the second equation is differential. Next, let us consider the boundary conditions that the functions from expansion (3) must satisfy.

For $z = 0$, the temperature and its derivative with respect to z , based on expansion (3), will have the form:

$$\begin{cases} T(0, t) = \zeta_0(t); \\ \left(\frac{\partial T(z, t)}{\partial z} \right)_{z=0} = \zeta_1(t) \end{cases} \quad (12)$$

Therefore, the third equation of the system of equations (1), taking into account (12), is transformed to the form:

$$\lambda \zeta_1(t) = -e\Theta \zeta_0^4(t) \quad (13)$$

This form completely coincides with equation (9) of Sedelnikov et al. (2024a).

For $z = h$, the temperature and its derivative with respect to z , based on decomposition (3), will have the form:

$$\begin{cases} T(h, t) = \zeta_0(t) + \zeta_1(t)h + \zeta_2(t)h^2 + \zeta_3(t)h^3; \\ \left(\frac{\partial T(z, t)}{\partial z} \right)_{z=h} = \zeta_1(t) + 2\zeta_2(t)h + 3\zeta_3(t)h^2. \end{cases} \quad (14)$$

Substituting (14) into the second equation of the system of equations (1) gives a result that coincides with equation (11) of Sedelnikov et al. (2024a) before linearisation:

$$\begin{aligned} \lambda [\zeta_1(t) + 2\zeta_2(t)h + 3\zeta_3(t)h^2] &= \\ &= Q - e\Theta [\zeta_0(t) + \zeta_1(t)h + \zeta_2(t)h^2 + \zeta_3(t)h^3]^4. \end{aligned} \quad (15)$$

Further in Sedelnikov et al. (2025), a linearisation with respect to h was performed. In this paper, the second-order terms with respect to h will be retained. After this simplification, (15) is transformed to the form:

$$\begin{aligned} \lambda [\zeta_1(t) + 2\zeta_2(t)h + 3\zeta_3(t)h^2] &= Q - \\ &- e\Theta \zeta_0^2(t) [\zeta_0^2(t) + 4\zeta_0(t)\zeta_1(t)h + 4\zeta_0(t)\zeta_2(t)h^2 + \zeta_1^2(t)h^2]. \end{aligned} \quad (16)$$

Equation (16) differs significantly from the linearised equation (12) in Sedelnikov et al. (2024a) and is a more accurate approximation of equation (15). We substitute conditions (11) in (16) and perform elementary transformations:

$$\begin{aligned} \frac{d\zeta_1(t)}{dt} &= -\frac{12ea\Theta}{\lambda} (a\beta_1 t^2 + T_0)^2 \zeta_1^2(t) - \frac{8ea\Theta}{\lambda h} (a\beta_1 t^2 + T_0)^3 \zeta_1(t) + \\ &+ \frac{2a}{\lambda h^2} [Q - 2\lambda\beta_1 th - 4e\Theta h^2 (a\beta_1 t^2 + T_0)^3 \beta_1 t]. \end{aligned} \quad (17)$$

Then finally, the functions in decomposition (3), taking into account the second approximation with respect to h , will take the form:

$$\begin{cases}
\zeta_0(t) = a\beta_1 t^2 + T_0; \\
\frac{d\zeta_1(t)}{dt} = -\frac{12ea\Theta}{\lambda} (a\beta_1 t^2 + T_0)^2 \zeta_1^2(t) - \frac{8ea\Theta}{\lambda h} (a\beta_1 t^2 + T_0)^3 \zeta_1(t) + \\
\quad + \frac{2a}{\lambda h^2} [Q - 2\lambda\beta_1 t h - 4e\Theta h^2 (a\beta_1 t^2 + T_0)^3 \beta_1 t]; \\
\zeta_2(t) = \beta_1 t; \\
\zeta_2(t) = -\frac{2e\Theta}{\lambda} (a\beta_1 t^2 + T_0)^2 \zeta_1^2(t) - \frac{4e\Theta}{3\lambda h} (a\beta_1 t^2 + T_0)^3 \zeta_1(t) + \\
\quad + \frac{1}{3\lambda h^2} [Q - 2\lambda\beta_1 t h - 4e\Theta h^2 (a\beta_1 t^2 + T_0)^3 \beta_1 t].
\end{cases} \quad (18)$$

The resulting system of functions (18) is a solution to the initial boundary value problem (1). The second equation (18) is solved only approximately.

4 Numerical modelling and comparison of results with a linearised model

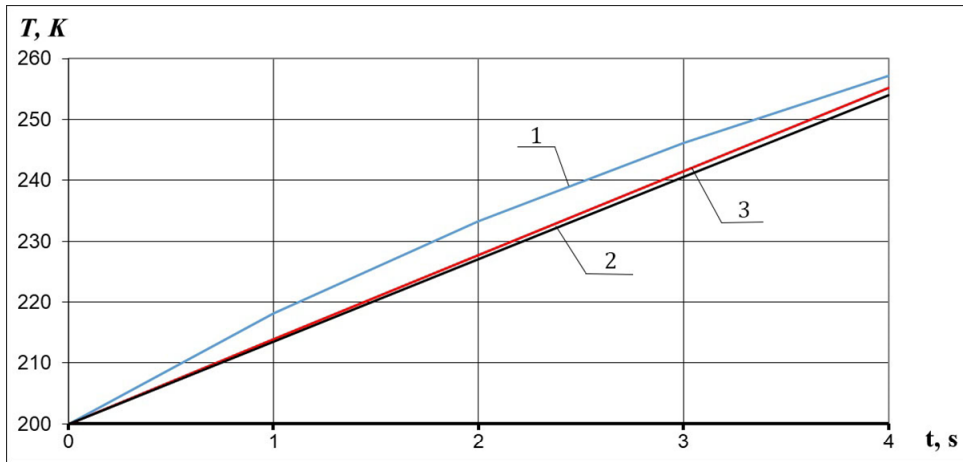
Let's perform numerical modelling in the Mathcad mathematical package. The values of the main characteristics of the plate and the temperature impact completely coincide with the work Sedelnikov et al. (2024a). The same values were used in Sedelnikov et al. (2021) for numerical simulation of temperature shock in the ANSYS package. This is done in order to make it possible to correctly compare the obtained simulation results. The main characteristics are shown in Table 1 (Sedelnikov et al., 2024a).

Table 1 The main parameters of the simulated plate

Parameter	Designation	Value	Dimension
Solar panel frame material	—	MA2	—
Coefficient of thermal conductivity	λ	96,3	1400
Stefan-Boltzmann constant	Θ		$W/(m^2 \cdot K^4)$
External heat flux	Q	1,400	W/m^2
Initial temperature of the solar panel frame	$T_0 = T(z, 0)$	200	K
Degree of blackness	e	0, 2	-
Specific heat	c	1,130.4	$J/(kg \cdot K)$
Density	ρ	1,780	g/m^3
Young's Module	E	$4 \cdot 10^{10}$	Pa
Shift modulus	μ		Pa
Poisson's Ratio	ν	0, 3	-
Solar panel length	l	1	m
Solar panel width	b	0, 5	m
Solar panel frame thickness	h	1	mm

Source: Sedelnikov et al. (2024a)

The following results were obtained in numerical simulation. They are shown in Figure 2.

Figure 2 Dynamics of surface layer temperature ($z = h$) (see online version for colours)

Notes: 1 – simulation results in ansys (Sedelnikov et al., 2021); 2 – simulation results based on the linearised model (Sedelnikov et al., 2024a); 3 – simulation results based on the proposed model, taking into account the terms of the second order relative to z .

It can be seen from Figure 2 that there are no significant differences from the linearised model (Sedelnikov et al., 2024a) taking into account the terms of the second order relative to z didn't give. Therefore, the use of a linearised model is fully justified for an approximate assessment of the effect of a thermal shock.

For the surface layer with the coordinate $z = 0$, the linearisation of equation (4) has no effect on the temperature distribution in the layer. This is a disadvantage not so much of the decomposition (2) as of the Fourier method (Watts, 2012). It is precisely such shortcomings that the authors of Arinchev (2022), as well as a number of other works, for example, Li et al. (2022) and Shenoy et al. (2010) point out.

5 Conclusions

The paper considers an approximate solution to the one-dimensional problem of thermal conductivity using expansion (3), taking into account the terms of the second order relative to z . The results obtained allow us to conclude that the considered approach did not bring a significant increase in accuracy compared to the linearised model in Sedelnikov et al. (2024a). Moreover, accounting for members of a higher order seems irrational. And this isn't due to the fact that the linearised model (Sedelnikov et al., 2024a) is accurate, but due to the disadvantages of the Fourier method itself in the context of the problem being solved. Apparently, the significant nonlinearity of the boundary conditions of the third kind makes the idea of using the decomposition of a function of two variables as a product of two independent functions of one variable less applicable. A similar position can be found in the works Belousova and Serdakova (2020) and Sedelnikov et al. (2023).

Thus, an important result of the work is to verify the reliability of the linearised model (Sedelnikov et al., 2024a). This model, in fact, has exhausted all the possibilities

of the Fourier decomposition method (Watts, 2012) and cannot be improved by taking into account a larger number of higher-order terms relative to z .

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Declarations

The authors declare that there is no conflict of interest.

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